



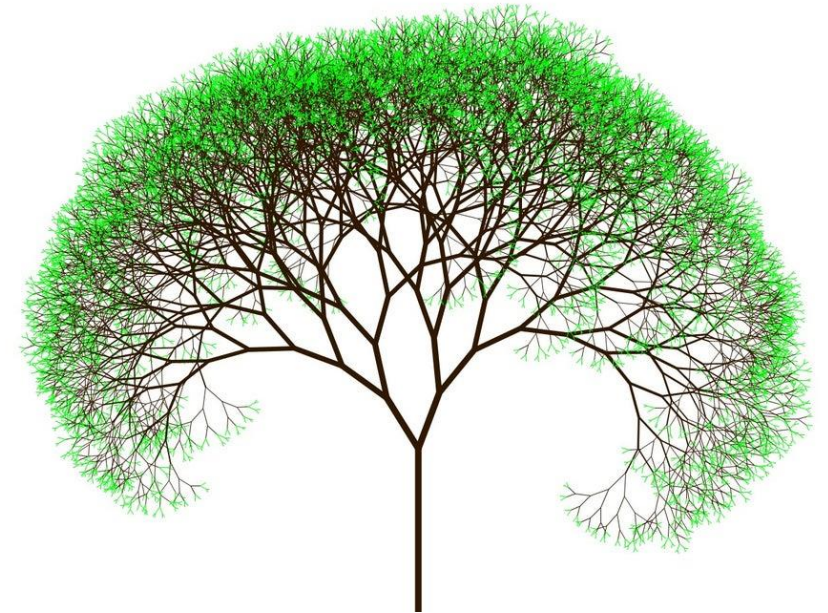
Institute for Research in
Fundamental Sciences

Proton structure functions at low x : the Fractal approach

Samira Shoeibi

in collaboration with

Shahin Atashbar Tehrani, Fatemeh Taghavi-Shahri



Weekly Seminar, School of Particles and Accelerators, 26 October 2022

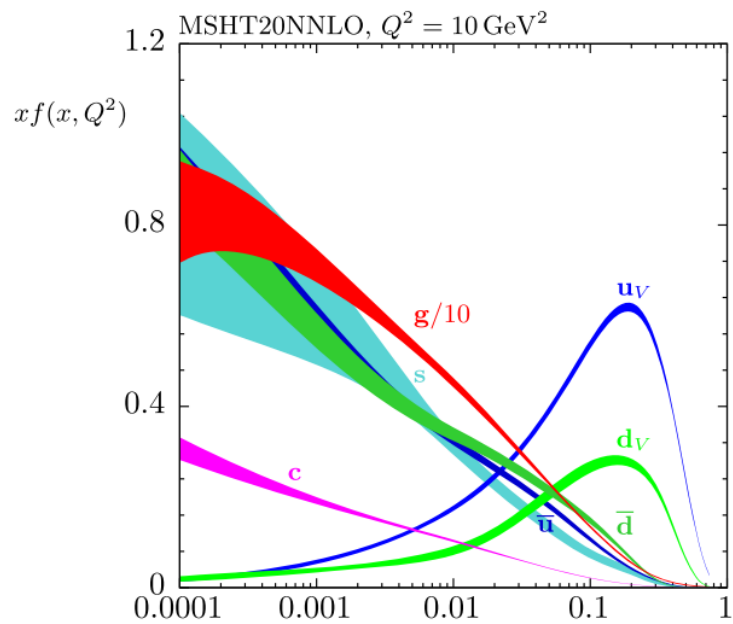
Outline

- Proton structure functions and PDFs
- Fractals
- Two essential ingredients to obtain PDFs:
 - Parametrizations of PDFs in initial value of Q_0^2
 - Experimental observable
- Minimization processes and Discussion of the QCD fit results
- Result

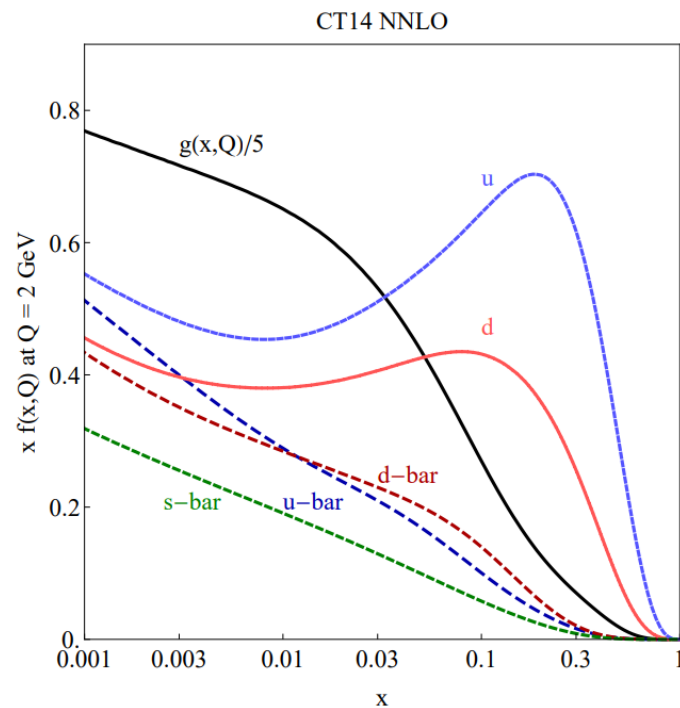
Proton structure functions and PDFs

- ❖ Quantum chromodynamics (**QCD**) describes the **properties and the internal structure of hadrons**.
- ❖ Hadrons are made up of **partons** (quarks and gluon) and at high energy scales (>200 MeV), these partons are probed almost free and the hadron's momentum are sharing among them. This information can be encoded in term of non-perturbative **Parton Distribution Functions (PDFs)**.
- ❖ PDFs are critical component in **High Energy Physics (HEP) phenomenology**, for example:
 - ✓ used in the computation of physical observables, such as cross sections.
 - ✓ are key component on the precision measurement of Standard Model parameters such as strong coupling constant.
- ❖ There are two main ways to **determining the PDFs**:
 - 1) using **lattice QCD**.
 - 2) using **large experimental data sets in global QCD analyses**
 - ✓ particular functional form for the x dependence of the PDFs (at a scale Q_0) like MSHT20 PDFs and CT14.
 - ✓ Neural Networks (NNs) like NNPDF.

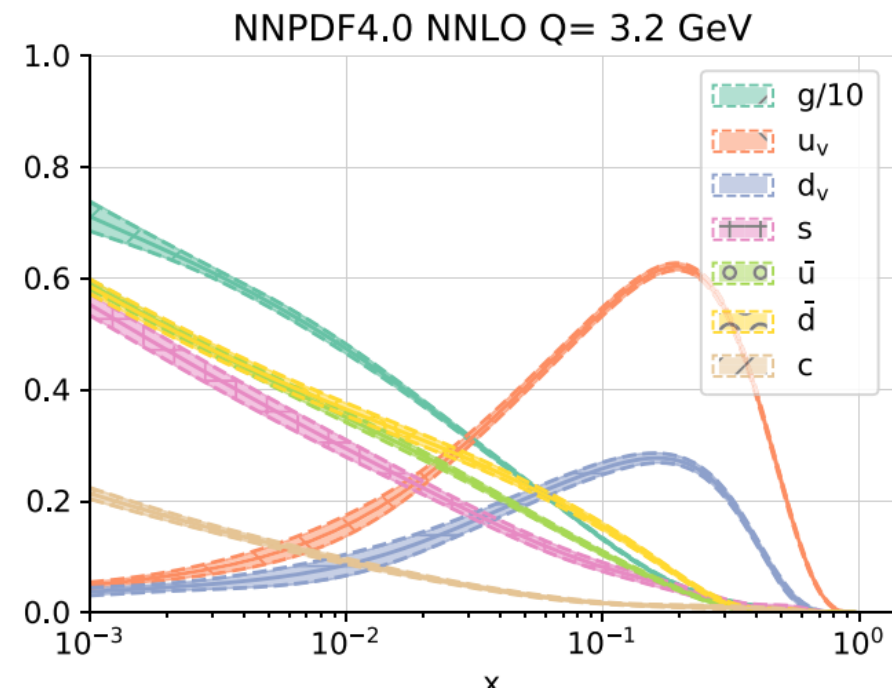
Proton structure functions and PDFs



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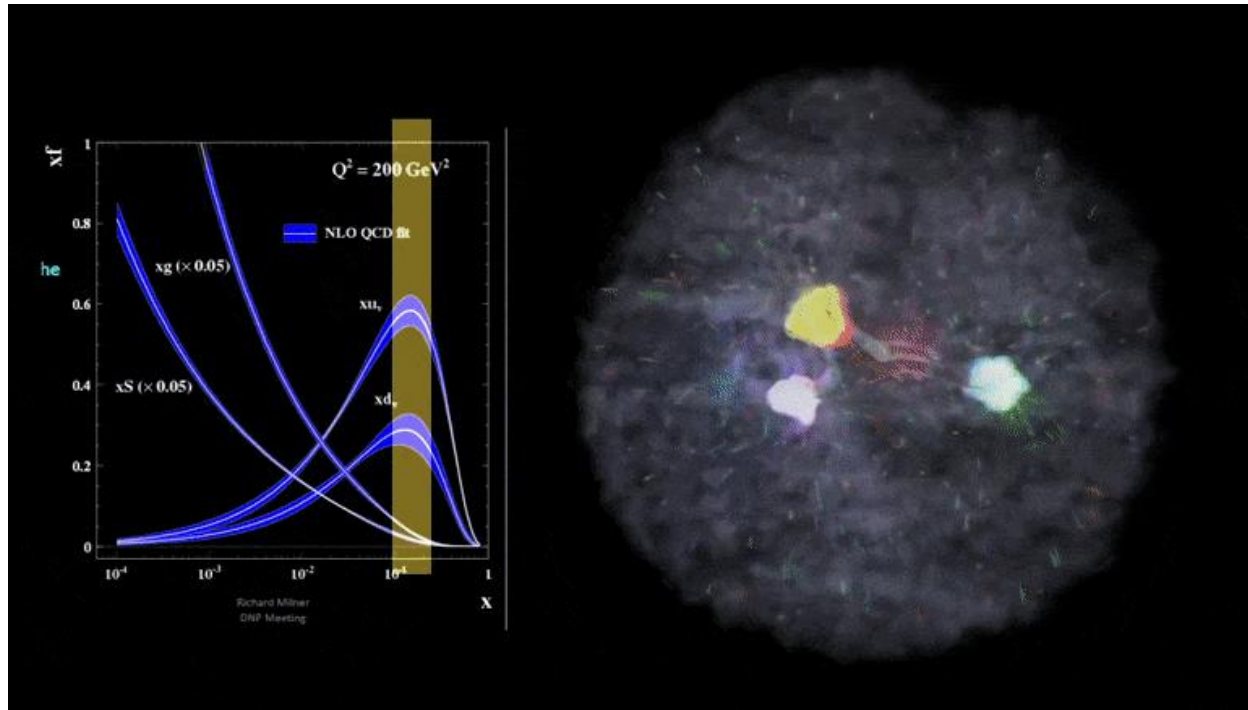


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Proton structure functions and PDFs



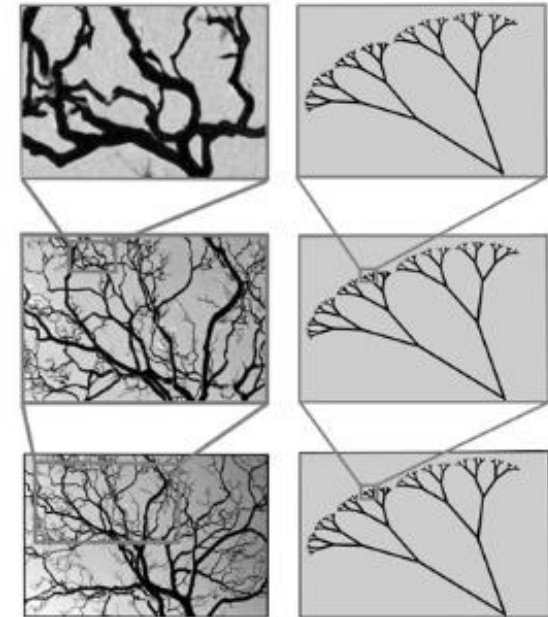
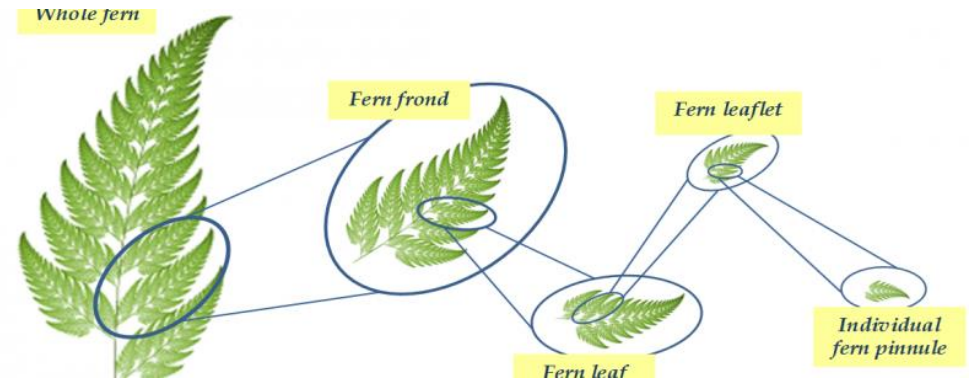
<https://scitechdaily.com/visualizing-the-proton-physicists-innovative-animation-depicts-the-subatomic-world-in-a-new-way/>

Fractals

Fractal is a curve or geometrical figure, **each part of which has the same statistical character as the whole** and they have some features such as:

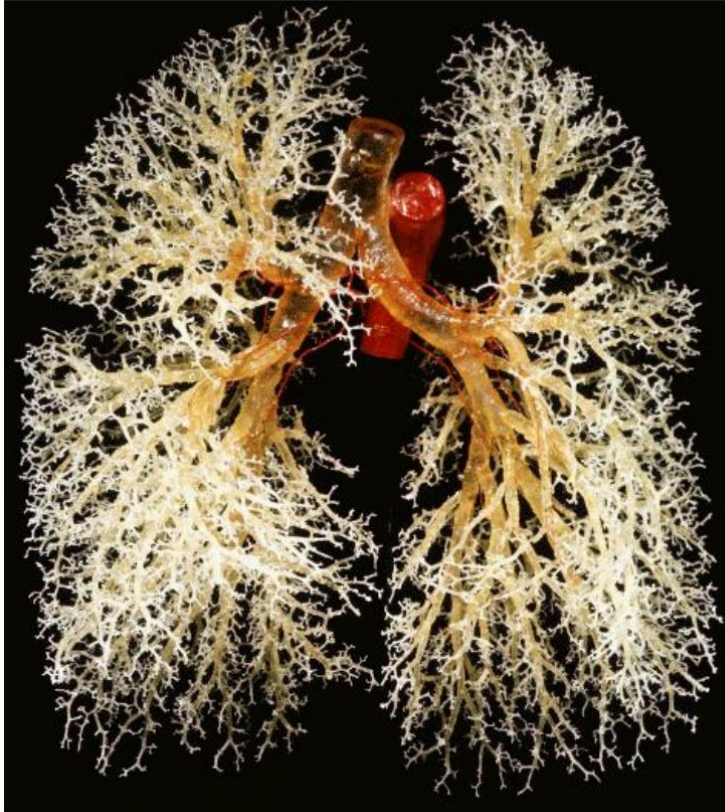
- **Self-similar**
- **Fractal Dimension**

Fractals in Nature:



Fractals

Fractals in Animal Bodies:

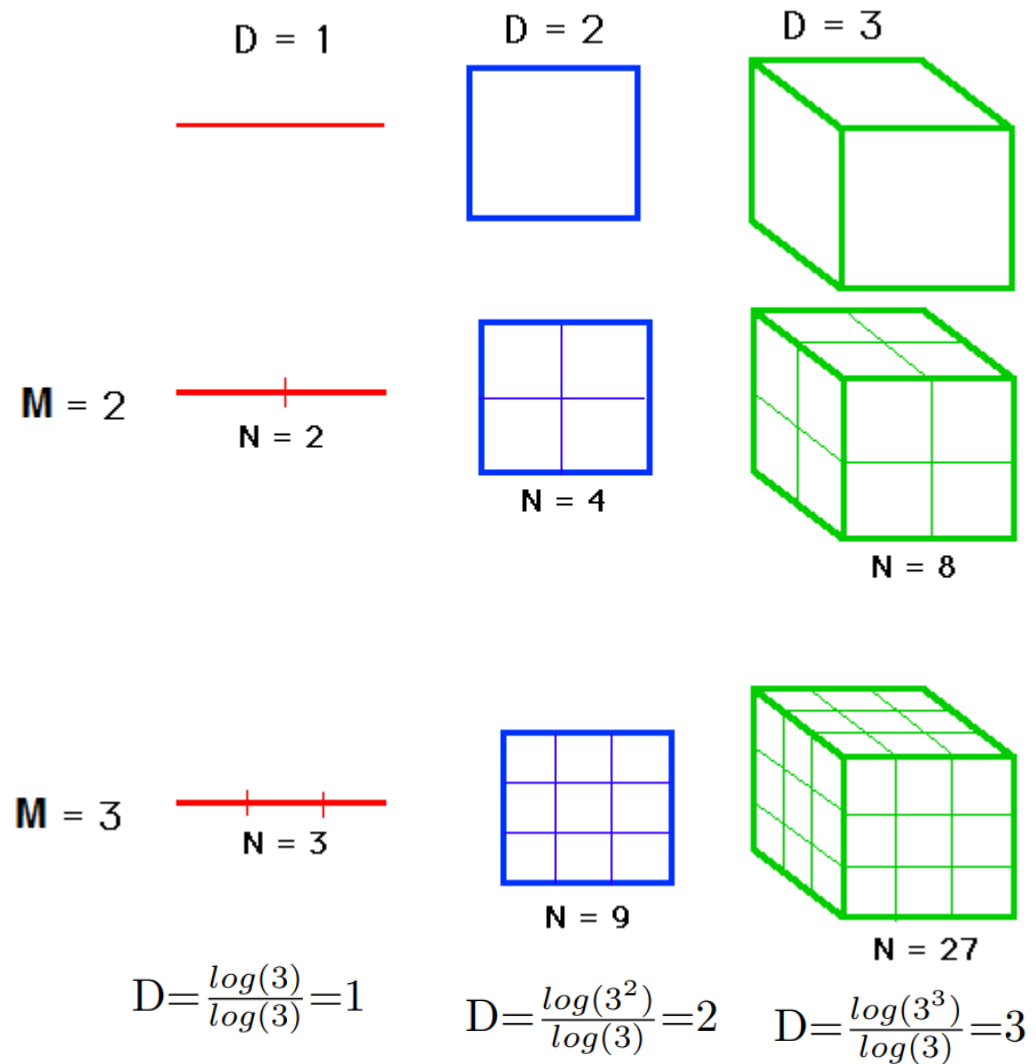


Fractal Lightning and Electricity:



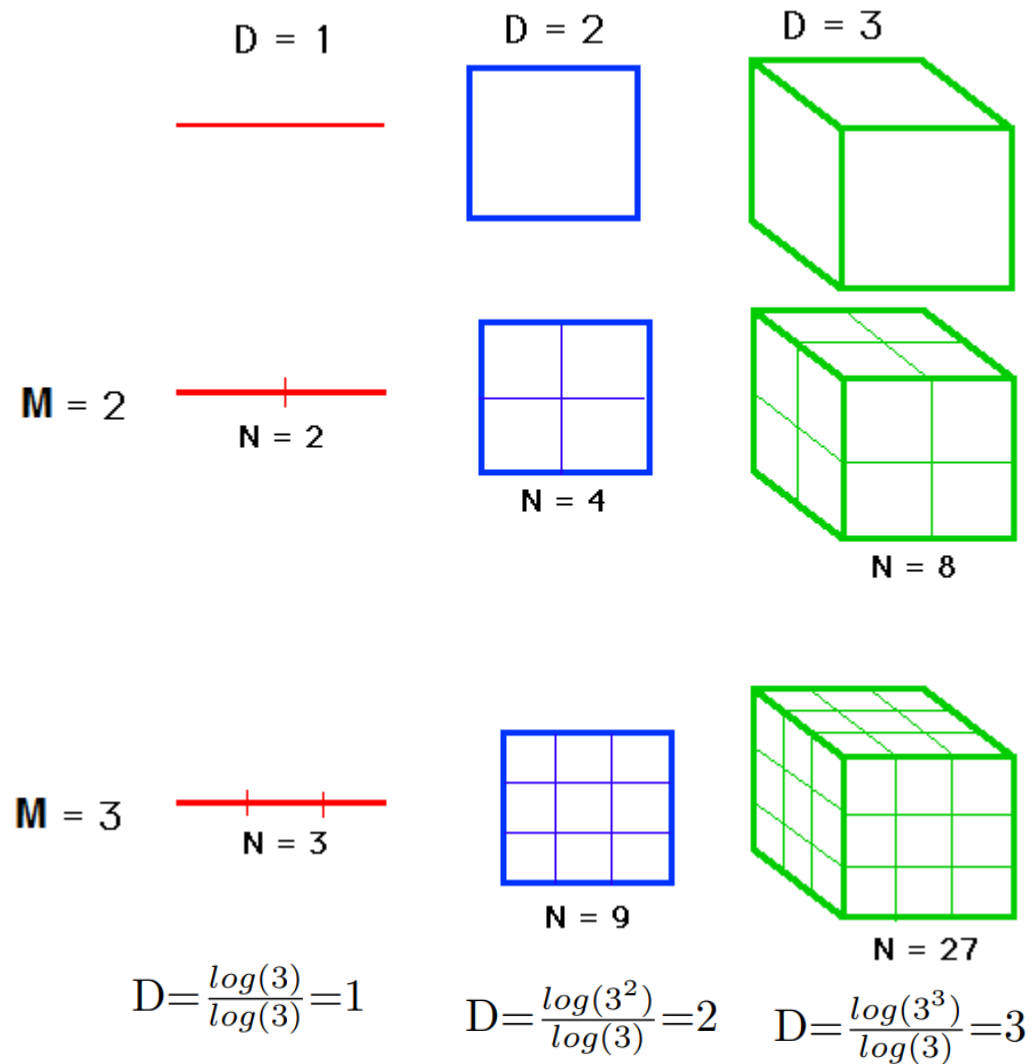
Fractals

Fractal Dimension:

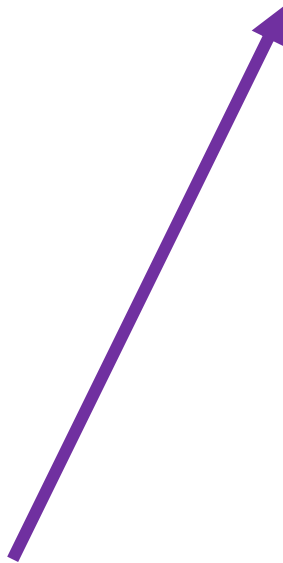


Fractals

Fractal Dimension:

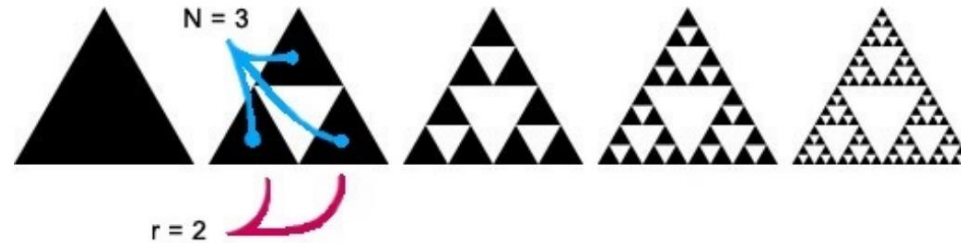


$$D = \frac{\log(\text{number of self similar objects})}{\log(\text{magnification factor})} = \frac{\log(N)}{\log(r)}$$



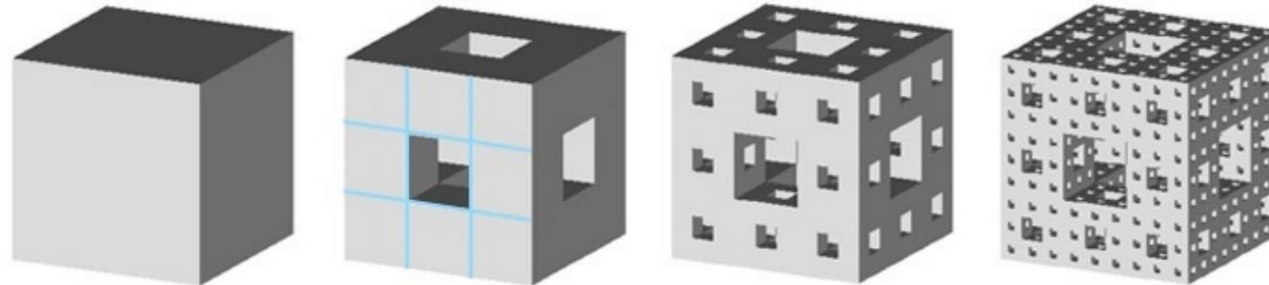
Fractals

Sierpinski Gasket:



$$D = \frac{\log(N)}{\log(r)} = \frac{\log(3)}{\log(2)} = 1.585$$

Menger Sponge:



$$D = \frac{\log(N)}{\log(r)} = \frac{\log(20)}{\log(3)} = 2.726$$

The fractal dimension describes how complicated or how large a self-similar object is. A plane is “larger” than a line. The Sierpinski gasket is not a line but also far from being a plane.

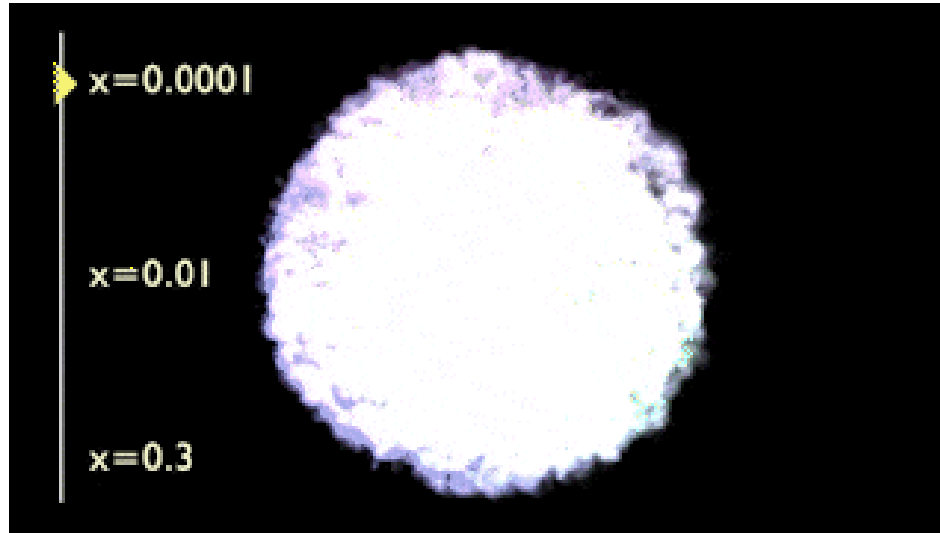
Fractals

If $M \equiv z$ and the number of the self-similar objects = $f(z)$, then

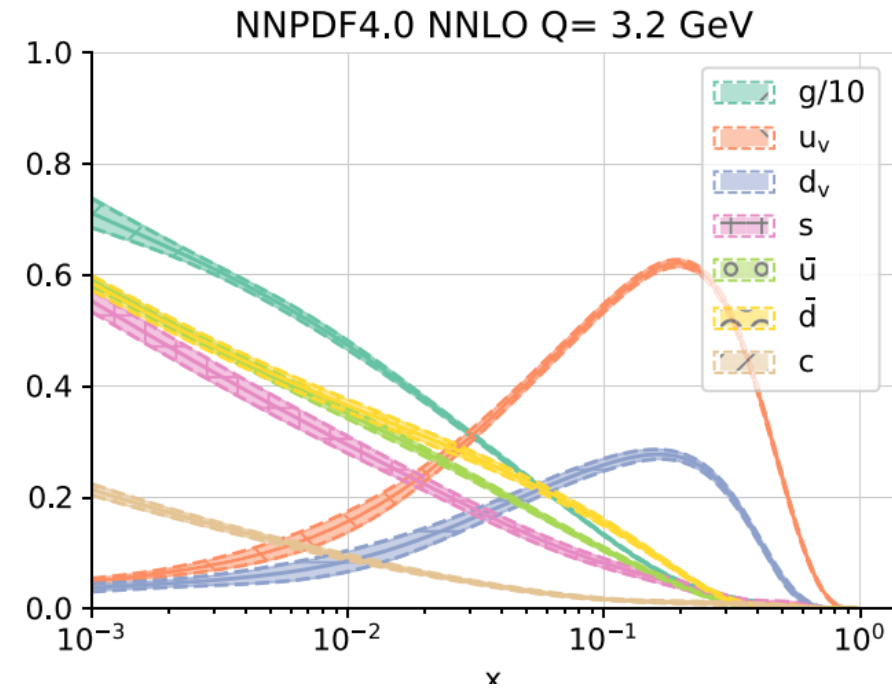
$$D = \frac{\partial \log f(z)}{\partial \log z} \implies \log f(z) = D \cdot \log z + D_0 \implies f(z) \propto z^D$$

In general, fractals may have two independent magnification factors, z and y . In this case the density $f(z, y)$ is written in the following way:

$$\log f(z, y) = D_{zy} \cdot \log z \cdot \log y + D_z \cdot \log z + D_y \cdot \log y + D_0$$



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- In **small x** ($x < 0.01$), **more gluon-gluon interactions** can be observed which increase the gluon and the sea quark densities. Therefore we can consider the proton as a soup of sea quarks and gluon and we can say in analogy to fractals, they follow self-similarity.

1) Parametrizations of PDFs in initial value of Q_0^2

Eur. Phys. J. C 24, 529–533 (2002)
Digital Object Identifier (DOI) 10.1007/s10052-002-0973-3

THE EUROPEAN
PHYSICAL JOURNAL C

Self-similar properties of the proton structure at low x

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² Charles University, Faculty of Mathematics and Physics, V Holešovičkách 2, 18000 Prague 8, Czech Republic

Received: 28 March 2002 /

Published online: 21 June 2002 – © Springer-Verlag / Società Italiana di Fisica 2002

Abstract. Self-similar properties of proton structure in the kinematic region of low values of the Bjorken variable x are introduced and studied numerically. A description of the proton structure function $F_2(x, Q^2)$ reflecting self-similarity is proposed with a few parameters which are fitted to recent HERA data. The specific parametrisation provides an excellent description of the data which cover the region of four momentum transfer squared $0.045 \leq Q^2 \leq 120 \text{ GeV}^2$, and of the Bjorken variable x $6.2 \cdot 10^{-7} \leq x \leq 0.01$.

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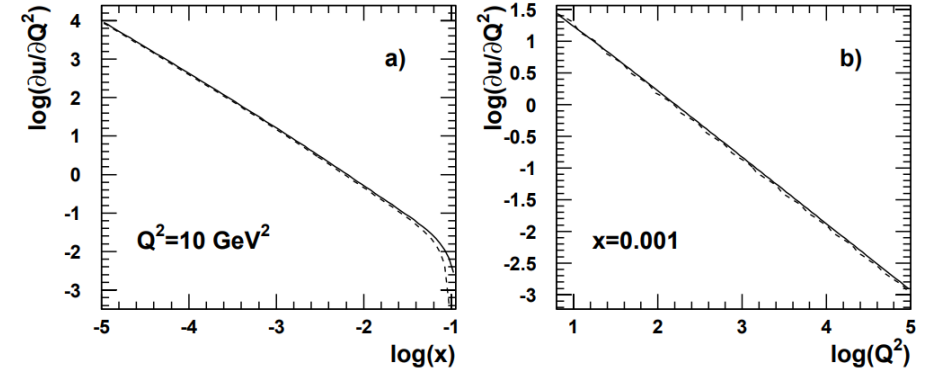


Fig. 2a,b. Logarithm of the unintegrated u -quark density $\partial u(x, Q^2) / \partial Q^2$ as a function of the Bjorken variable x **a** and Q^2 **b**. The full and dashed lines correspond to GRV parametrizations in LO and NLO [10], respectively

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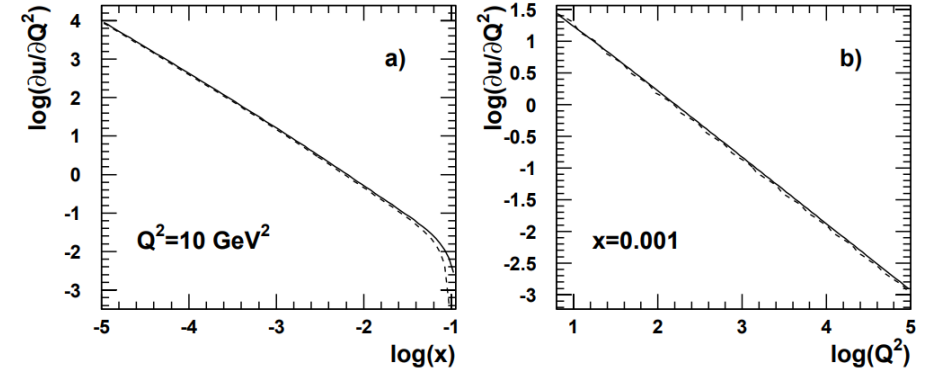


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$$\log(f_{q/p}(x, k_t^2)) = D_1^q \log\left(\frac{1}{x}\right) \log\left(1 + \frac{k_t^2}{q_0^2}\right) + D_2^q \log\left(\frac{1}{x}\right) + D_3^q \log\left(1 + \frac{k_t^2}{q_0^2}\right) + D_0^q$$

Quark Parton Model \longrightarrow $F_2 = x \sum_i e_i^2 (q_i + \bar{q}_i)$

1) Parametrizations of PDFs in initial value of Q_0^2

Towards precision determination of uPDFs

Magnus Hansson¹ and Hannes Jung²

1- Lund University

2- DESY, FRG

The unintegrated Parton Density Function of the gluon is obtained from a fit to dijet production in DIS as measured at HERA. Reasonable descriptions of the measurements are obtained, and a first attempt to constrain the intrinsic transverse momentum distribution at small k_{\perp} is presented [1].

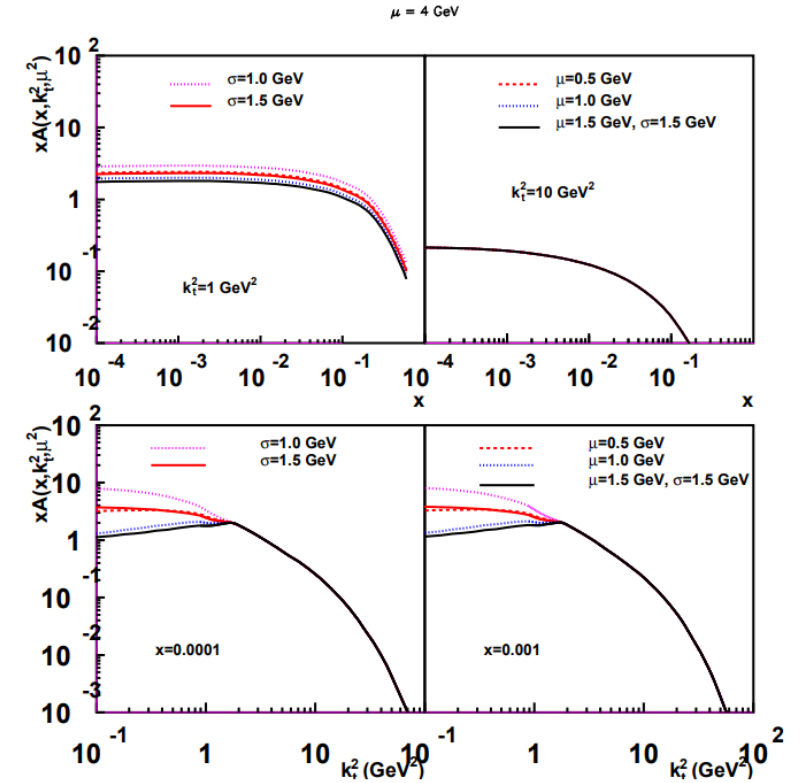


Figure 1: The unintegrated gluon distribution at a scale $\bar{q} = 4$ GeV for different values of μ and σ of the intrinsic k_{\perp} distribution as a function of x for fixed k_{\perp} (top) and as a function of k_{\perp} (bottom) for fixed x

1) Parametrizations of PDFs in initial value of Q_0^2

$$\log(f_{q/p}(x, k_t^2)) = D_1^q \log\left(\frac{1}{x}\right) \log\left(1 + \frac{k_t^2}{q_0^2}\right) + D_2^q \log\left(\frac{1}{x}\right) + D_3^q \log\left(1 + \frac{k_t^2}{q_0^2}\right) + D_0^q - \log(M^2)$$

$$x f_{g/p}(x, k_t^2) = \left(\frac{1}{x}\right)^{B^g} (1-x)^{C^g} (1-D^g x) \cdot e^{-\frac{(\mu-k_t)^2}{\sigma^2}} \frac{1}{M^2}$$

$$x f_{i/p}(x, Q^2) = \int_0^{Q^2} x f_{i/p}(x, k_t^2) dk_t^2$$

**Initial
PDFs**

$$\left\{ \begin{array}{l} x f_{q/p}(x, Q_0^2) = \frac{e^{D_0^q q_0^2} x^{-D_2^q + 1}}{M^2 (1 + D_3^q - D_1^q \log(x))} \left(x^{-D_1^q \log\left(1 + \frac{Q_0^2}{q_0^2}\right)} \left(1 + \frac{Q_0^2}{q_0^2}\right)^{D_3^q + 1} - 1 \right) \\ x f_{g/p}(x, Q_0^2) = A' \left(\frac{1}{x}\right)^{B^g} (1-x)^{C^g} (1-D^g x) \end{array} \right.$$

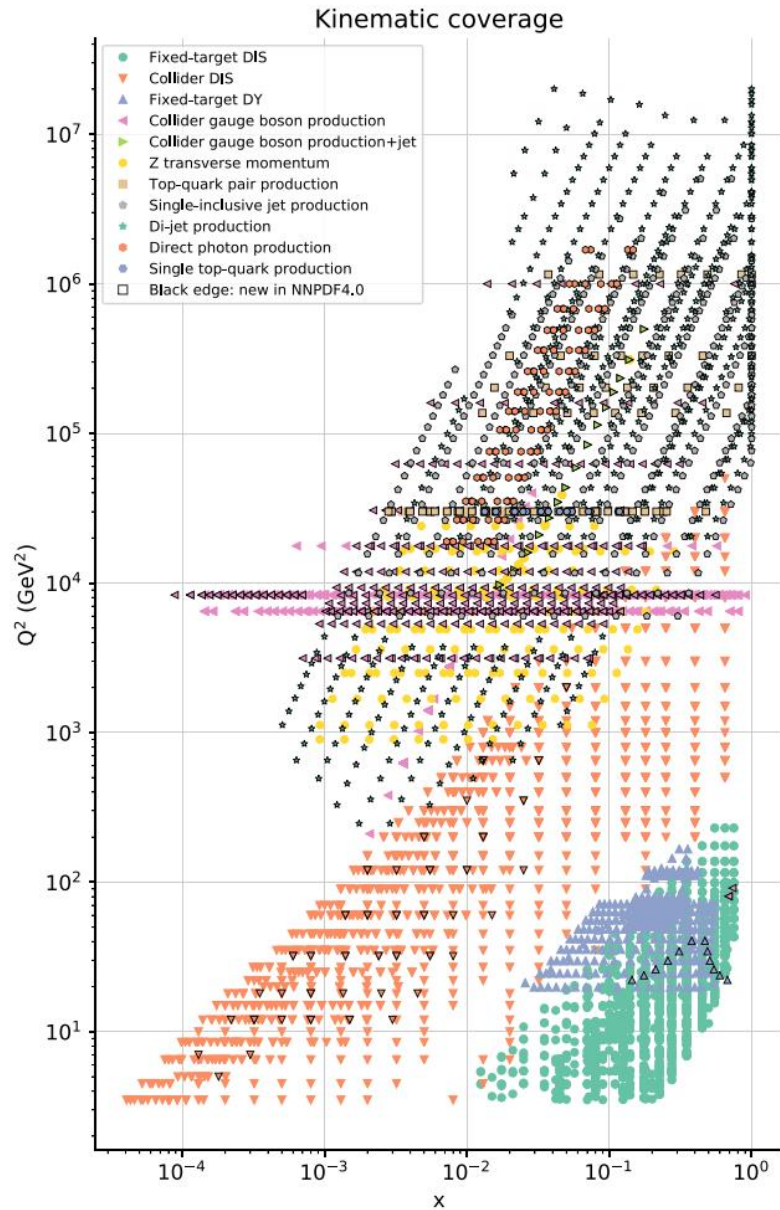
here A' is:

$$A' = \frac{A}{M^2} \left(\exp^{-\frac{(\mu-1)^2}{\sigma^2}} + \exp^{-\frac{\mu^2}{\sigma^2}} \right) \sqrt{\pi} \mu \left(\text{Erf}\left(\frac{1-\mu}{\sigma}\right) + \text{Erf}\left(\frac{\mu}{\sigma}\right) \right)$$

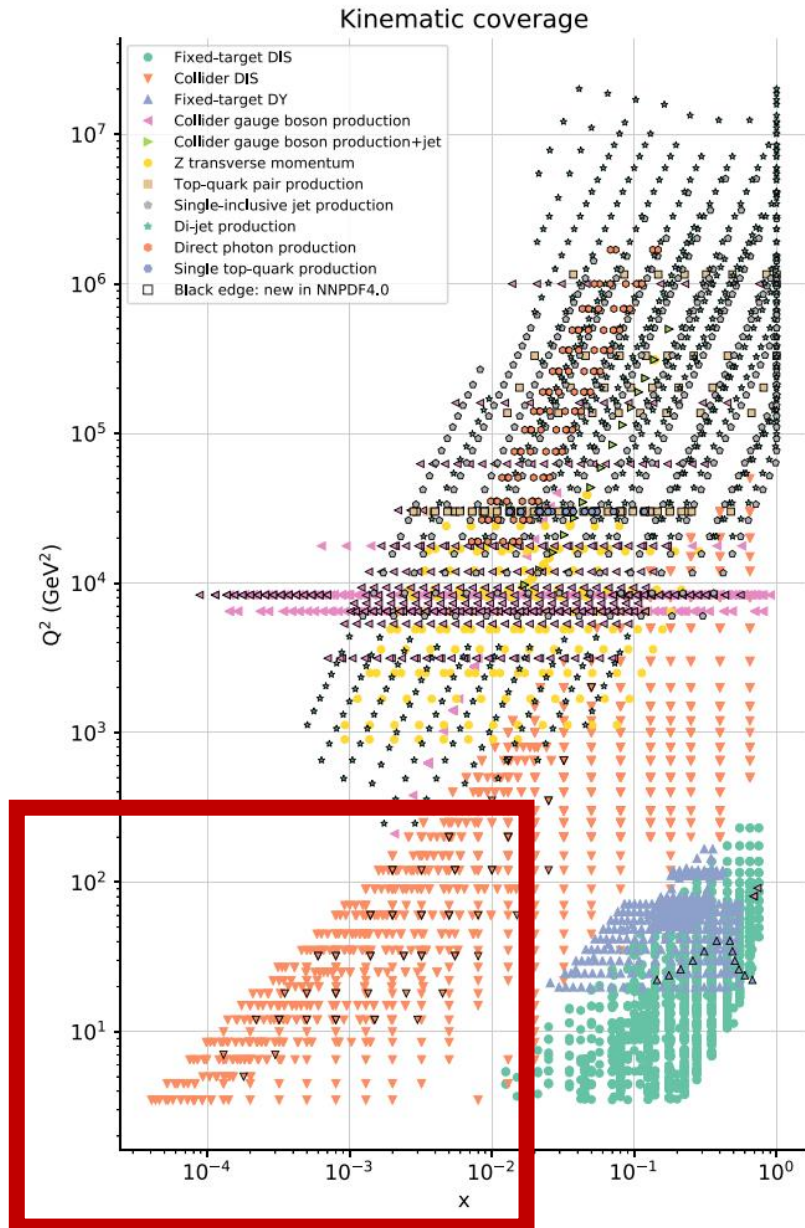
where $\text{Erf}(z)$ is the "error function" which is defined by

$$\text{Erf}(z) \equiv \frac{2}{\pi} \int_0^z \exp^{-t^2} dt$$

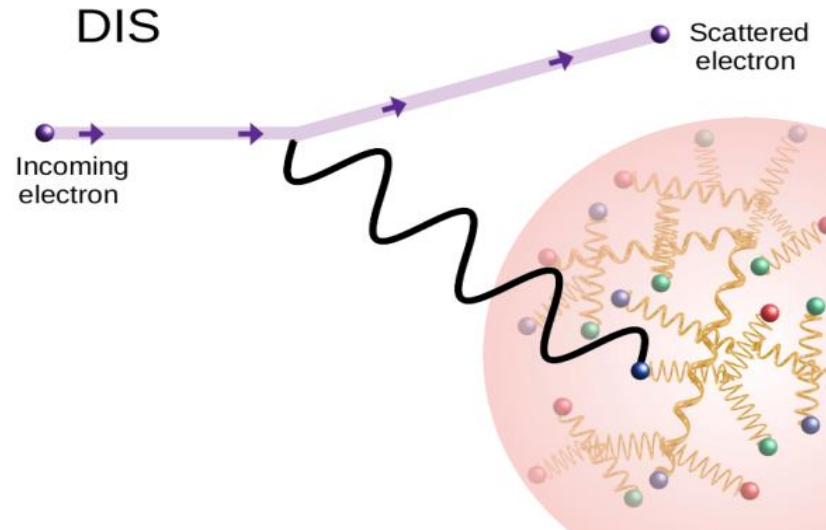
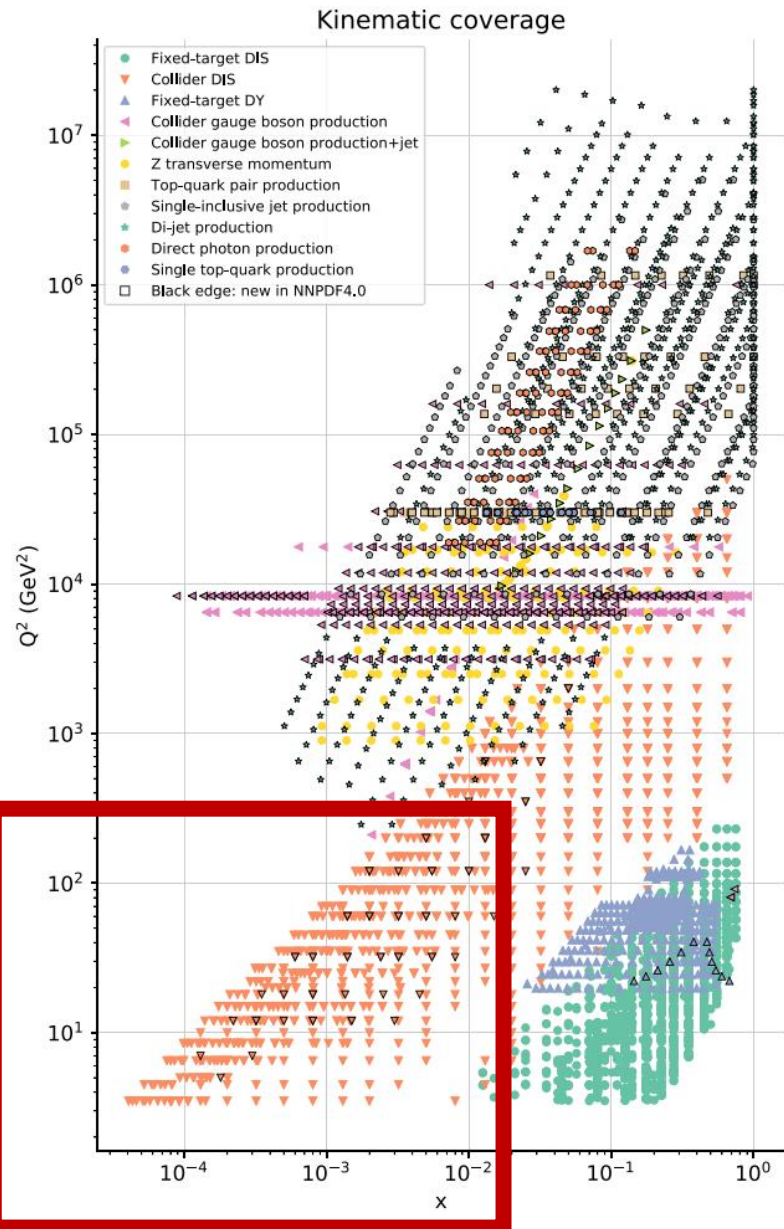
2) Experimental observable



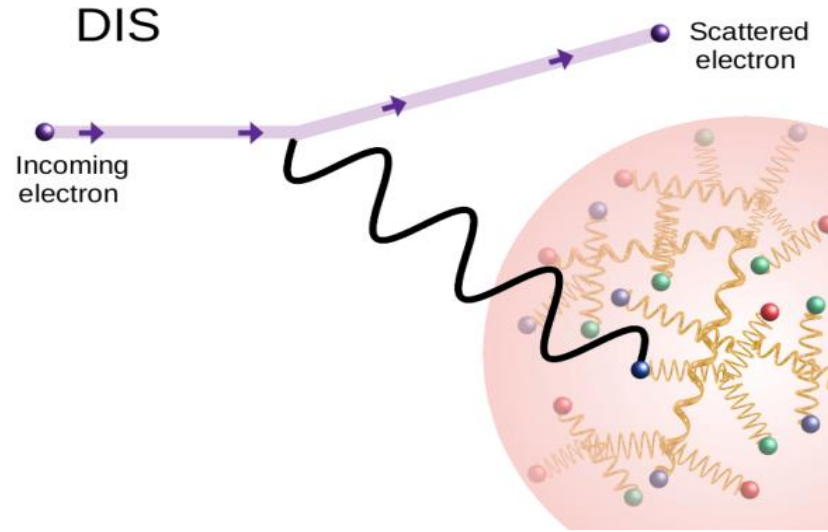
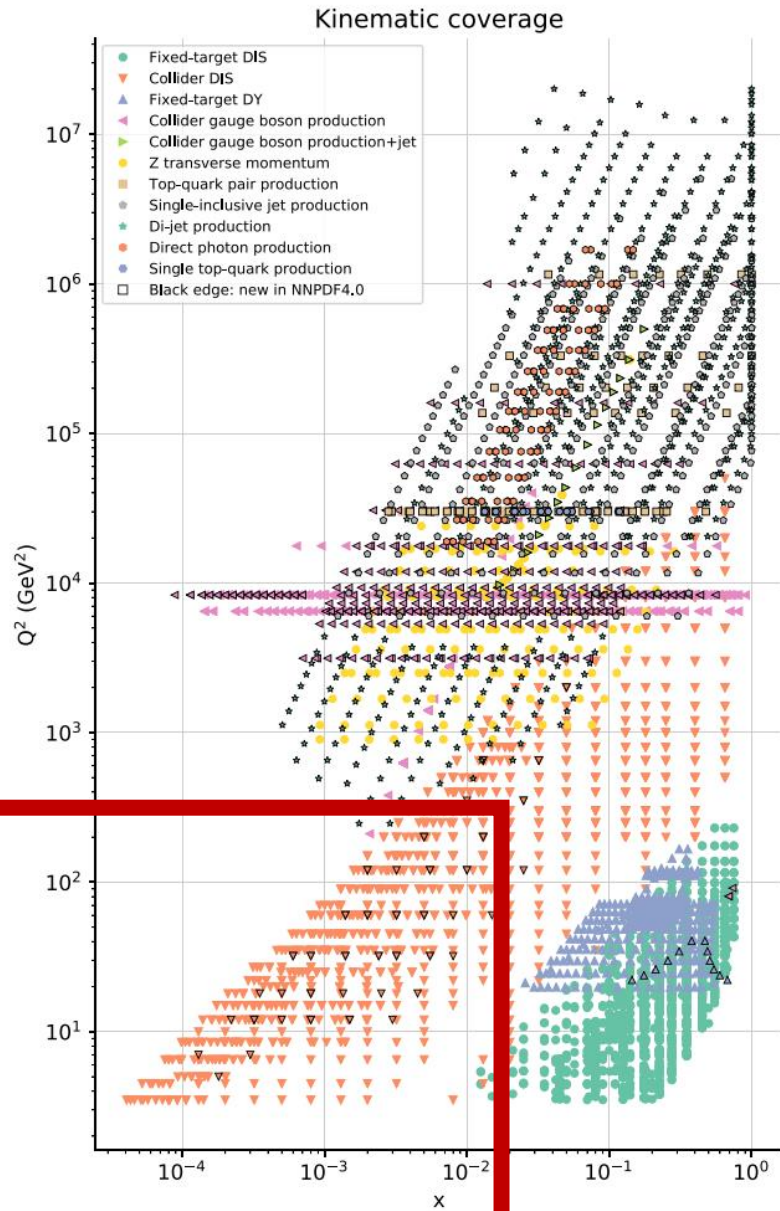
2) Experimental observable



2) Experimental observable



2) Experimental observable



Q^2 GeV ²	x_{Bj}	$\sigma_{r,NC}^+$	δ_{stat} %	δ_{uncor} %	δ_{cor} %	δ_{rel} %	$\delta_{\gamma p}$ %	δ_{had} %	δ_1 %	δ_2 %	δ_3 %	δ_4 %	δ_{tot} %
15	0.200×10^{-3}	1.256	3.21	3.61	1.00	0.54	-0.11	-0.14	-0.01	-0.01	-0.48	-0.15	4.99
15	0.246×10^{-3}	1.360	1.17	1.52	1.43	0.76	-1.11	-0.46	-0.13	0.01	-0.39	-0.10	2.82
15	0.320×10^{-3}	1.290	0.94	1.36	1.00	0.68	-0.63	-0.22	-0.05	0.00	-0.19	-0.10	2.17
18	0.268×10^{-3}	1.318	3.26	3.67	0.99	0.54	-0.11	-0.17	-0.01	-0.01	-0.46	-0.15	5.06
18	0.328×10^{-3}	1.362	1.25	1.67	1.09	0.62	-0.67	-0.47	0.06	0.00	-0.18	-0.08	2.58
18	0.500×10^{-3}	1.266	0.93	1.01	0.95	0.71	-0.44	-0.29	-0.08	0.00	-0.14	-0.07	1.89
22	0.500×10^{-3}	1.335	2.57	1.22	1.18	1.36	-0.85	-0.43	-0.08	0.02	-0.07	0.04	3.50
27	0.335×10^{-3}	1.389	4.11	3.92	0.98	0.54	-0.12	-0.15	-0.01	-0.01	-0.45	-0.14	5.81
27	0.410×10^{-3}	1.405	1.16	2.10	1.15	0.63	-0.70	-0.40	-0.04	-0.01	-0.08	-0.11	2.85
27	0.500×10^{-3}	1.407	1.15	1.46	0.92	0.77	-0.59	-0.31	-0.02	0.00	-0.21	-0.07	2.32
27	0.800×10^{-3}	1.286	2.46	0.83	0.87	0.77	-0.48	-0.28	0.02	-0.02	0.02	0.04	2.89
35	0.574×10^{-3}	1.507	1.36	2.01	0.99	0.59	-0.47	-0.35	-0.03	-0.03	-0.19	-0.14	2.76
35	0.800×10^{-3}	1.399	0.95	0.81	0.91	0.95	-0.60	-0.25	-0.01	-0.01	-0.15	-0.17	1.94
45	0.800×10^{-3}	1.486	1.81	1.45	1.44	0.80	-1.20	-0.53	-0.06	-0.07	0.08	-0.09	3.14
45	0.130×10^{-2}	1.306	0.05	0.62	0.86	0.73	-0.03	0.05	0.00	-0.02	-0.07	-0.21	1.61

Minimization processes and Discussion of the QCD fit results

[H1](#) and [ZEUS](#)

Collaborations
experimental data

(Eur.Phys.J.C 75 (2015) 12, 580)

Parametrized form of
PDFs

Minimization processes and Discussion of the QCD fit results

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Collaborations
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(*Eur.Phys.J.C* 75 (2015) 12, 580)

Parametrized form of
PDFs



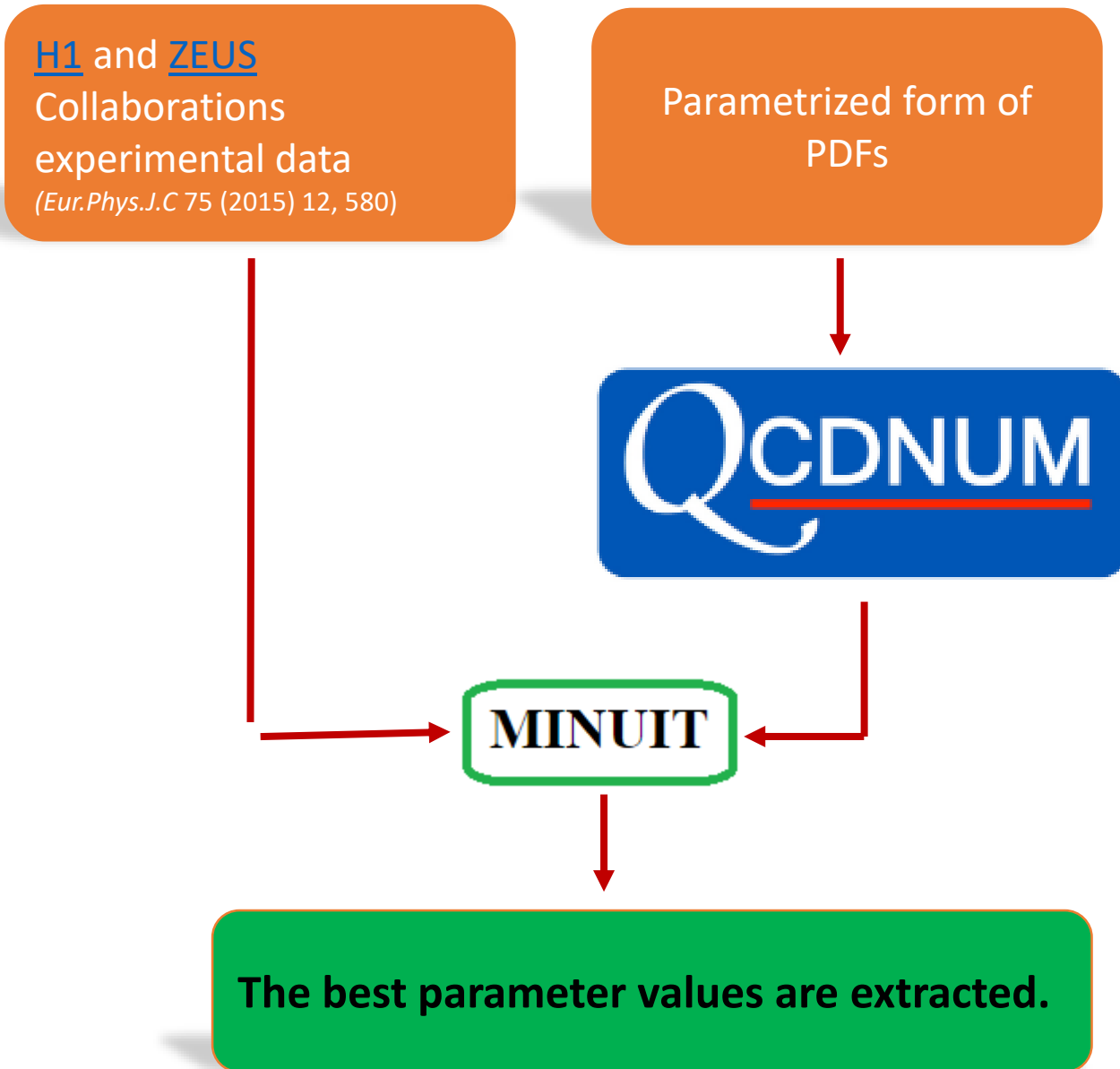
$$\frac{dq_i(x, Q^2)}{d \ln Q^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{d\xi}{\xi} \left[\sum_j q_j(\xi, Q^2) P_{qq}^{(0)}\left(\frac{x}{\xi}\right) + g(\xi, Q^2) P_{qg}^{(0)}\left(\frac{x}{\xi}\right) \right]$$

$$\frac{dg(x, Q^2)}{d \ln Q^2} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{d\xi}{\xi} \left[\sum_j q_j(\xi, Q^2) P_{gq}^{(0)}\left(\frac{x}{\xi}\right) + g(\xi, Q^2) P_{gg}^{(0)}\left(\frac{x}{\xi}\right) \right]$$

$$F_2(x, Q^2) = x \sum_{i=q, \bar{q}, g} \int_x^1 \frac{d\xi}{\xi} C_2^i\left(\frac{x}{\xi}, Q^2, \mu_f^2, \mu_r^2\right) f_i(\xi, \mu_f^2)$$

$$\sigma_{r, NC}^{\pm} = F_2(x, Q^2) - \frac{y}{Y_+} F_L(x, Q^2) .$$

Minimization processes and Discussion of the QCD fit results



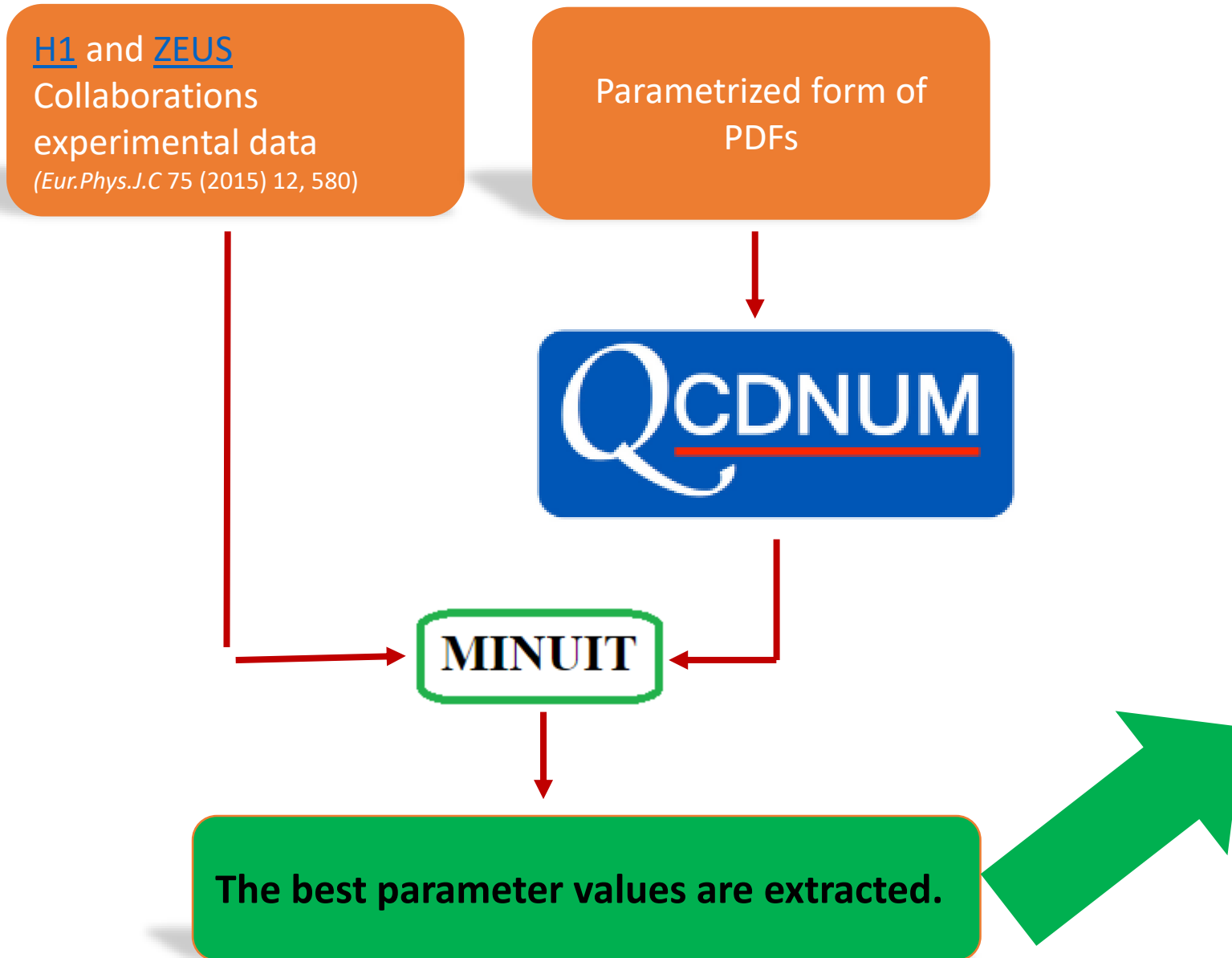
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$$\sigma_{r, NC}^{\pm} = F_2(x, Q^2) - \frac{y}{Y_+} F_L(x, Q^2) .$$

Minimization processes and Discussion of the QCD fit results



Parameters	$xf_{q/p}(x, Q_0^2)$
D_0^q	4.81 ± 0.01
q_0	0.0187 ± 0.0001
D_1^q	-0.0051 ± 0.0009
D_2^q	1.138 ± 0.002
D_3^q	-1.285 ± 0.005

Parameters	$xf_{g/p}(x, Q_0^2)$
A	0.191 ± 0.005
B^g	-3.01 ± 0.03
C^g	4^*
D^g	-3071.1 ± 80.6
σ	7.15 ± 6.18
μ	1.5^*

Minimization processes and Discussion of the QCD fit results

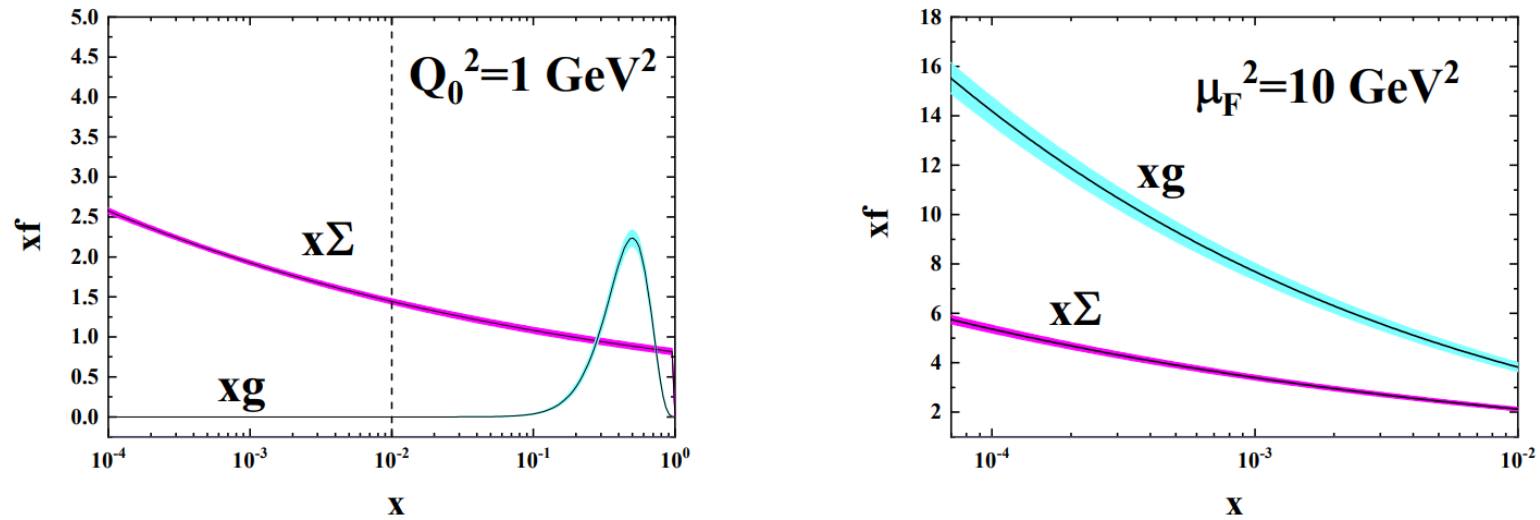


Fig. 3. (left) The parton distribution functions at $Q_0^2 = 1 \text{ GeV}^2$. (Right) The same but for $Q^2 = 10 \text{ GeV}^2$ and in the low x region below $x < 0.01$. Note that the fractal PDFs are valid only for $x < 0.01$.

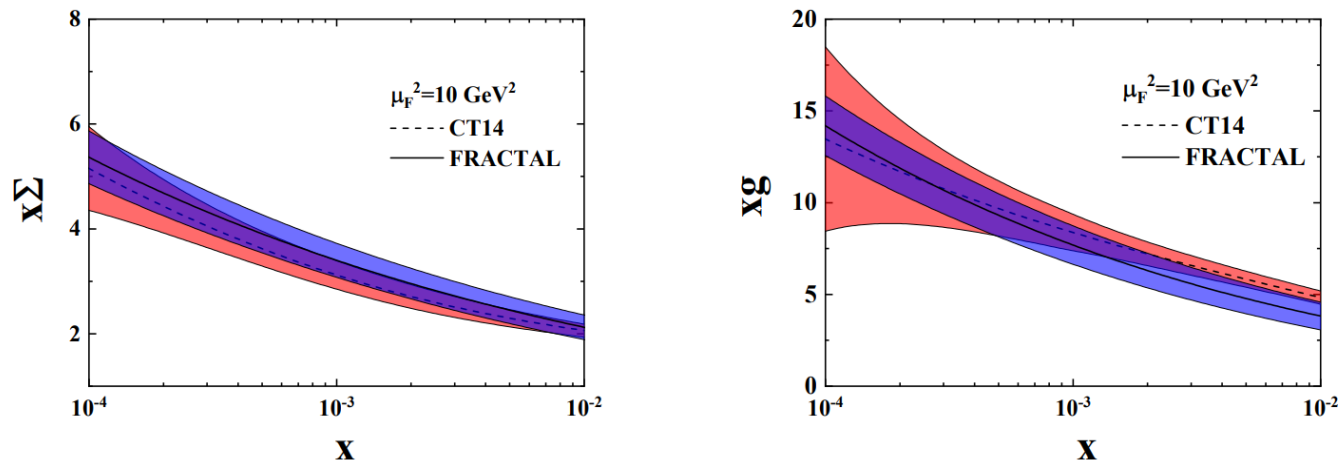


Fig. 5. The parton distribution functions with uncertainties band at low x in comparison with those from CT14 at $Q^2 = 10 \text{ GeV}^2$.

Minimization processes and Discussion of the QCD fit results

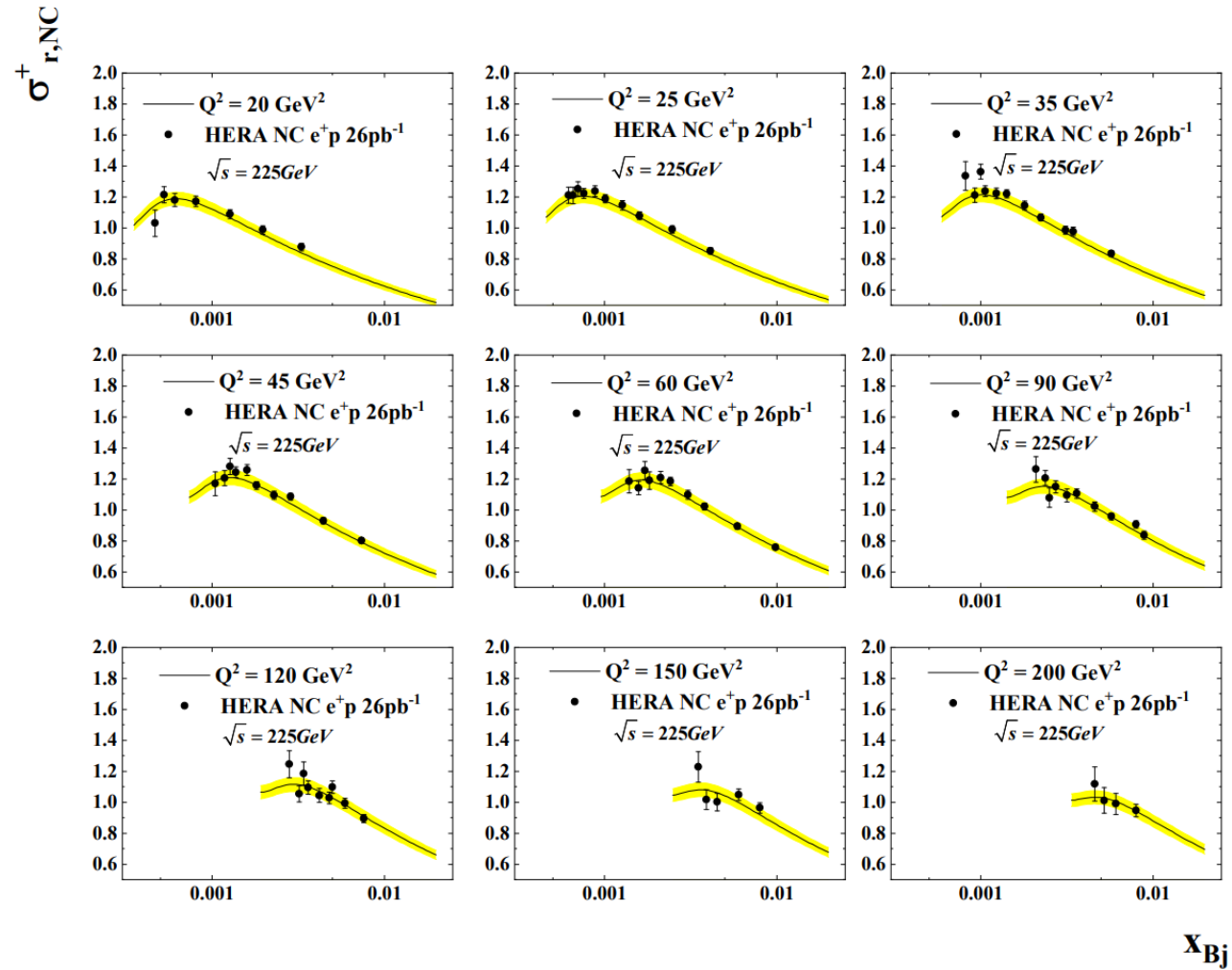


Fig. 6. The prediction for the reduced cross-section in the low x region below $x < 0.01$ for $Q^2 = 20\text{GeV}^2$ to $Q^2 = 200\text{GeV}^2$ and $\sqrt{s} = 225\text{GeV}$. Data points are from NC interactions in HERA positron-proton DIS processes.

Minimization processes and Discussion of the QCD fit results

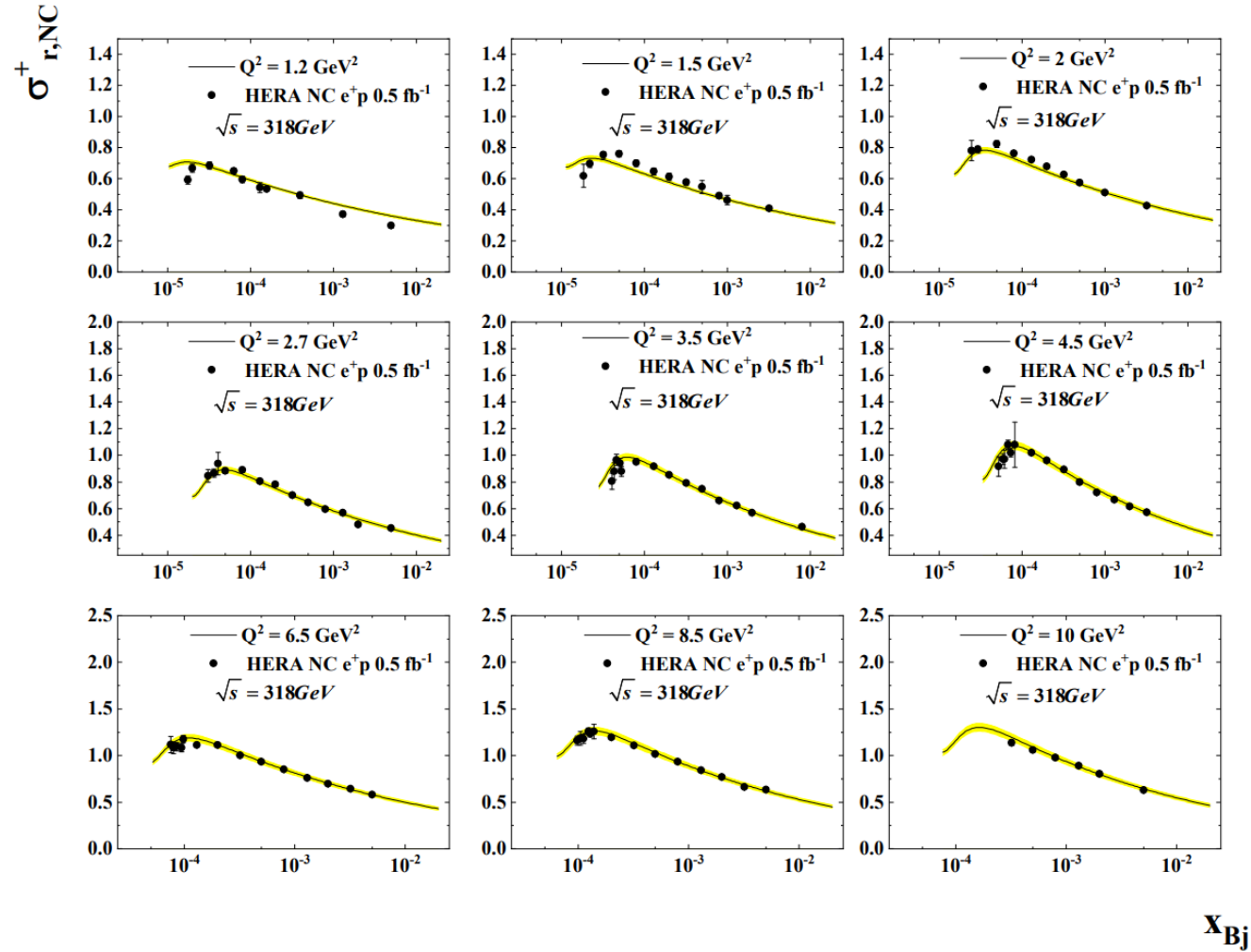


Fig. 7. The prediction for the reduced cross-section in the low x region below $x < 0.01$ for $Q^2 = 1.2 \text{ GeV}^2$ to $Q^2 = 10 \text{ GeV}^2$ and $\sqrt{s} = 318 \text{ GeV}$. Data points are from NC interactions in HERA positron-proton DIS processes.

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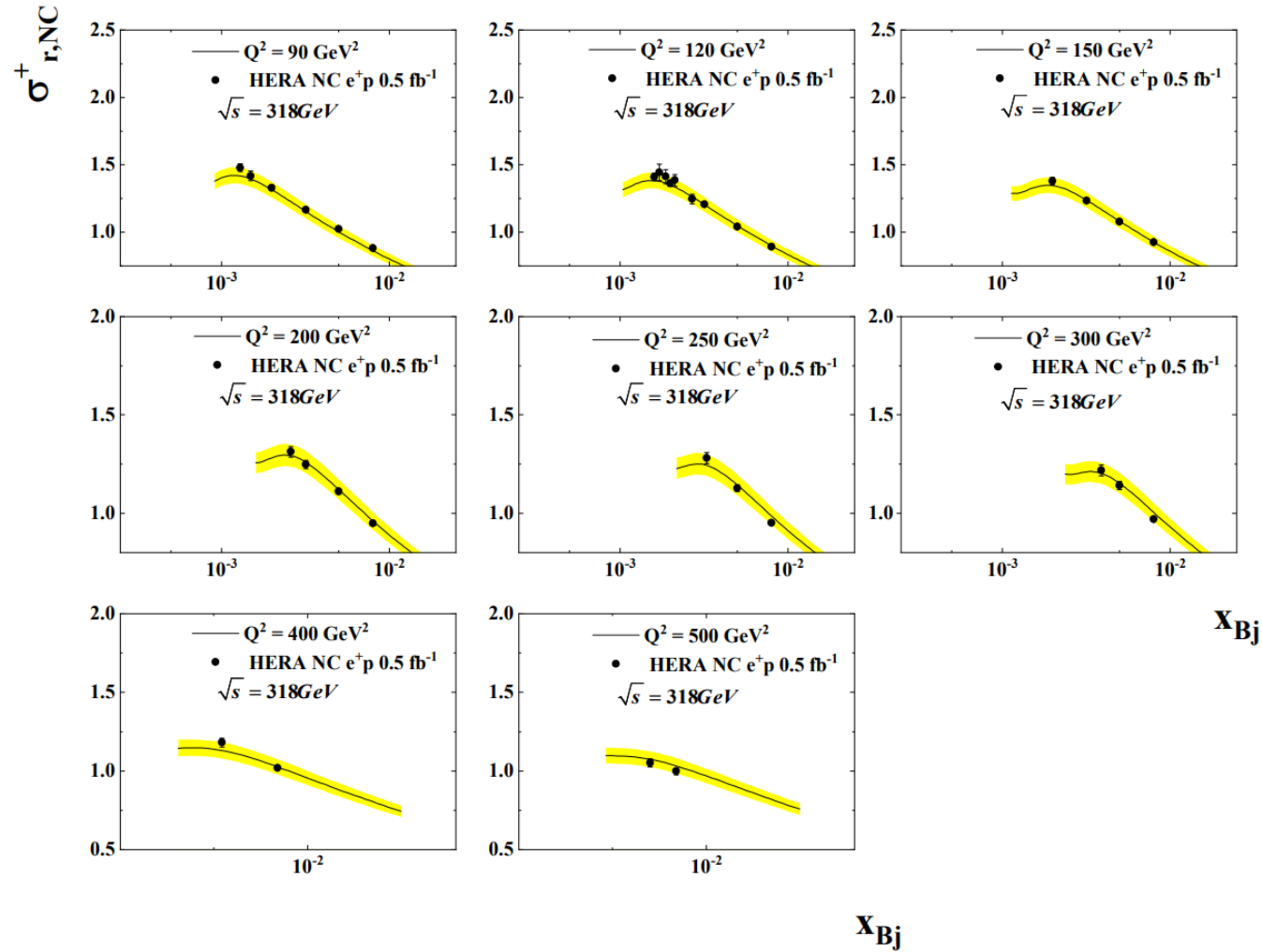


Fig. 9. The prediction for the reduced cross-section in the low x region below $x < 0.01$ for $Q^2 = 90 \text{ GeV}^2$ to $Q^2 = 500 \text{ GeV}^2$ and $\sqrt{s} = 318 \text{ GeV}$. Data points are from NC interactions in HERA positron-proton DIS processes.

Result

- ❖ **Fractal PDFs can describe low x physics well.**



Back up

$$\chi_{\text{global}}^2(\{\xi_i\}) = \sum_{m=1}^{m^{\text{exp}}} w_m \chi_m^2(\{\xi_i\}),$$

$$\chi_m^2(\{\xi_i\}) = \left(\frac{1 - \mathcal{N}_m}{\Delta \mathcal{N}_m} \right)^2 + \sum_{j=1}^{N_m^{\text{data}}} \left(\frac{(\mathcal{N}_m \mathcal{O}_j^{\text{data}} - \mathcal{T}_j^{\text{theory}}(\{\xi_i\}))}{\mathcal{N}_m \Delta_j^{\text{data}}} \right)^2.$$

$$\frac{\chi^2}{d.o.f} = 1.148$$